

Airfoil Design (continued)

2) Inverse Design Techniques

* Pressure distribution is prescribed over the airfoil and airfoil shape is directly computed.

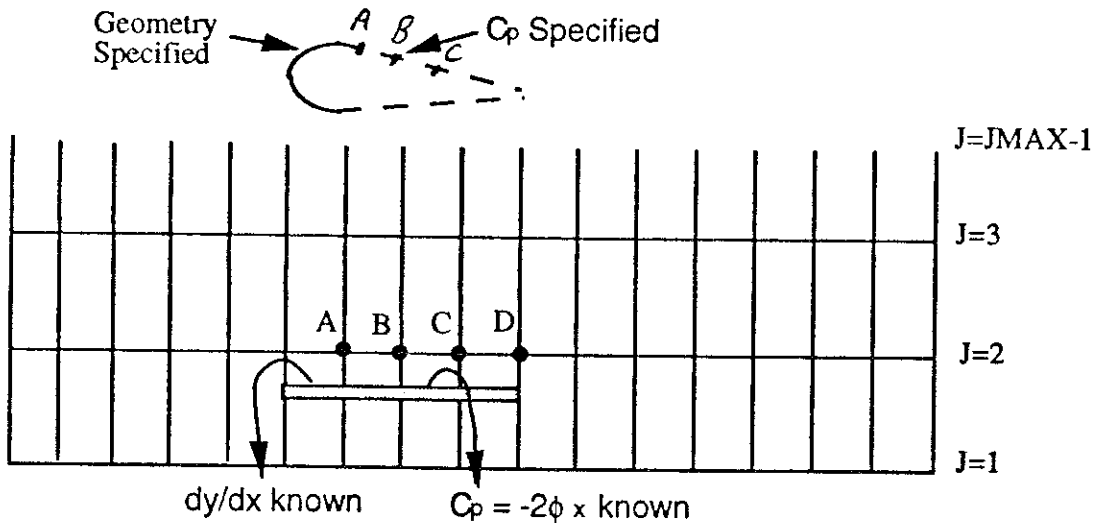
- Unlike optimization techniques, there is no objective function to be minimized.

• Several inverse design techniques exist. See web notes (Chapter 6) for references. One such method is...

Carlson's Inverse Design Technique

- May be used with TSD on a Cartesian grid.
- May also be used with FPE on a curvilinear grid (body fitted)
- For simplicity, let's consider here application to TSD in a non-lifting flow...

In Carlson's technique, only a part of the airfoil geometry is specified near the nose. Over the rest of the surface, the surface pressure distribution is prescribed.



Using Carlson's approach, we next construct a Cartesian grid and the TSD equation is solved using the principles learnt earlier in this course. At all points, including the points at $j=2$, away from the wall, the TSD equation is solved as:

$$(1 - \mu_{ij}) \frac{P_{i+\frac{1}{2},j} - P_{i-\frac{1}{2},j}}{\Delta x_{i+1} + \Delta x_i} + (\mu_{ij}) \frac{P_{i-\frac{1}{2},j} - P_{i-\frac{3}{2},j}}{\Delta x_{i+1} + \Delta x_i} + \frac{Q_{i,j+\frac{1}{2}} - Q_{i,j-\frac{1}{2}}}{\Delta y_{j+1} + \Delta y_j} = 0$$

Murman & Cole

At points adjacent to the airfoil, the above equation is modified depending on whether the point has the body-slope prescribed at $i, j-1/2$ or c_p prescribed. If the body slope is prescribed, then the discretized equation is

$$(1 - \mu_{ij}) \frac{P_{i+\frac{1}{2},j} - P_{i-\frac{1}{2},j}}{\Delta x_{i+1} + \Delta x_i} + (\mu_{ij}) \frac{P_{i-\frac{1}{2},j} - P_{i-\frac{3}{2},j}}{\Delta x_{i+1} + \Delta x_i} + \frac{Q_{i,j+\frac{1}{2}} - \left. \frac{dy}{dx} \right|_i}{\Delta y_{j+1} + \Delta y_j} = 0$$

Note: $Q_{i,j-\frac{1}{2}} = \left. \frac{dy}{dx} \right|_{j-\frac{1}{2}} = V_\infty \left(\frac{dy}{dx} \right)_{body} = \left(\frac{dy}{dx} \right)_{body}$

At a point such as B, where only the c_p distribution is known, ϕ is prescribed from the relationship

$$\phi_B = \phi_A + \int_A^B \underbrace{(\phi_x)}_{\phi_x} dx$$

where:

$\phi \equiv$ Perturbation Velocity Potential

Pressure Coefficient:

$$C_p \equiv \frac{P - P_0}{\frac{1}{2} \rho_0 V_0^2} = \frac{P - P_0}{\frac{\gamma}{2} P_0 M_0^2}$$

Dynamic Pressure

(show, at home, that these are equal)

$$C_p = \frac{P/P_0 - 1}{\frac{\gamma}{2} M_0^2}$$

Recall our isentropic relation:

$$P = P_0 \left[1 - \frac{\gamma-1}{2} M_0^2 \left(\frac{V^2}{V_0^2} - 1 \right) \right]^{\frac{\gamma}{\gamma-1}}$$

Combine these last two expressions (and see next page) to show that...

$$\text{For small disturbances: } C_p \approx -\frac{2\phi_x}{V_0}$$

Thus the ϕ_B equation above becomes

$$\phi_B = \phi_A + \int_A^B \left(-\frac{V_0 C_p}{2} \right) dx \Rightarrow \phi_B = \phi_A - \frac{V_0}{2} \int_A^B C_p dx$$

If non-dimensionalize so $V_0 = 1$, then $\phi_B = \phi_A - \frac{1}{2} \int_A^B C_p dx$

3.6 Small Disturbance Approximation for Surface pressure Coefficient, C_p :

We next develop a small disturbance approximation for the surface pressure. Our starting point is equations (3.2) and (3.4), which may be combined to yield:

$$\begin{aligned}
 C_p &= \frac{\left[1 + \frac{\gamma-1}{2} M_\infty^2 \left(1 - \frac{u^2 + v^2}{V_\infty^2}\right)\right]^{\frac{\gamma}{\gamma-1}} - 1}{\frac{\gamma}{2} M_\infty^2} \\
 &= \frac{\left[1 + \frac{\gamma-1}{2} M_\infty^2 \left(1 - \frac{(V_\infty + \phi_x)^2 + \phi_y^2}{V_\infty^2}\right)\right]^{\frac{\gamma}{\gamma-1}} - 1}{\frac{\gamma}{2} M_\infty^2} \\
 &= \frac{\left[1 + \frac{\gamma-1}{2} M_\infty^2 \left(-2 \frac{\phi_x}{V_\infty}\right)\right]^{\frac{\gamma}{\gamma-1}} - 1}{\frac{\gamma}{2} M_\infty^2}
 \end{aligned}$$

In arriving at the above form, we have neglected second powers of terms such as ϕ_x/V_∞ and ϕ_y/V_∞ .

Next, we use the binomial expansion

$$(1 + \epsilon)^n \approx 1 + n\epsilon \tag{3.20}$$

where ϵ is any small quantity $\epsilon \ll 1$ and n is any real number. Then,

$$C_p = \frac{\left[1 - \frac{\gamma}{2} M_\infty^2 \frac{\phi_x}{V_\infty} - 1\right]}{\frac{\gamma}{2} M_\infty^2} \xrightarrow{\text{Correction}} C_p \approx \frac{\left[1 + \frac{\gamma}{2} M_\infty^2 \left(-2 \frac{\phi_x}{V_\infty}\right)\right] - 1}{\frac{\gamma}{2} M_\infty^2}$$

$$\boxed{C_p = -2 \frac{\phi_x}{V_\infty}} \quad \Leftarrow \text{For small disturbance flow}$$

*

where A is the aftmost point of the airfoil where slope is prescribed. Since ϕ is explicitly known at point B, and all the downstream points C, D, etc., the governing equation is not solved here. The SLOR procedure along the y-lines containing points B, C, D, etc. is applied between J=3 and JMAX-1.

After each relaxation sweep, ϕ_z is known from the latest sweep. Therefore, $\phi_B, \phi_C,$ etc. may be computed. The relaxation process is repeated several hundred times until ϕ everywhere ceases to change.

After the relaxation process has converged, dy/dx values ^{are} found at the solid surface immediately beneath the points B, C, D, etc. from the ϕ field using the relationship

*

$$\frac{dy}{dx} = \phi_y \Big|_{i, j+\frac{1}{2}}$$

Conversely, these values may be found from the TSD equation rewritten as

$$\frac{dy}{dx} \Big|_i = Q_{i, j+\frac{1}{2}} + \frac{\Delta y_{j+1} + \Delta y_j}{\Delta x_{i+1} + \Delta x_i} \left[(1 - \mu_{ij}) (P_{i+\frac{1}{2}, j} - P_{i-\frac{1}{2}, j}) + \mu_{ij} (P_{i-\frac{1}{2}, j} - P_{i-\frac{3}{2}, j}) \right]$$

Knowing dy/dx everywhere, the airfoil is computed as

$$y(x) = y_A + \int_A^x \left(\frac{dy}{dx} \right) dx$$

Carlson's technique is extremely simple. It may be incorporated into any TSD or FPE program. It may easily be extended to a 3-D wing design as well.

In Carlson's approach, when the final airfoil shape is obtained, the trailing edge occasionally remains open or has a cross-over creating a fish-tail. The nose region must be opened up more, or pressed down, and the design performed again until the trailing edge closes properly.

